TDDFT Studies of Electronic Structure and First Hyperpolarizability of Tetra-nuclear Cubane-like Transition Metal Clusters

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Abstract: The TDDFT method is first applied in a series of tetra-nuclear transition metal clusters studies for nonlinear optical properties. The results indicate that the charge transfer inside the metal core $[MCu_3X_4]$ (M=W, Mo; X=S, O, Cl, Se, Br) makes contribution to the optical nonlinearity. It is possible to enhance the hyperpolarizability by substituting the ligands of the clusters.

Keywords: Nonlinear optical properties, transition metal clusters, TD- DFT method.

A series of tri-nuclear transition metal clusters with incomplete cubane-like configurations have been studied by TDDFT method. The calculations show that they have enormously large second-order polarizabilities (β) and are potential nonlinear optical materials for infrared double frequency conversion¹. In this paper some tetra-nuclear transition metal clusters with cubane-like configurations, MCu₃X₄ (PPh₃)₃ (M=W, Mo; X= S, O, Cl, Se, Br), were studied by TDDFT method for a reference. The structure-NLO property relationship was discussed for future NLO material design.

The structures of model clusters of the study were shown in **Figure 1**. All the structural parameters were taken from the X-ray structural data². The PPh₃ ligands were simplified by simpler (PH₃) ligands to enhance the computational efficiency. The first-order hyperpolarizabilities, β , of these molecular clusters were calculated by TDDFT method at GGA/LB94 level³. The electronic structures were calculated at GGA/BLYP level⁴. The basis sets for transition metal atoms and nonmetals were triple- ζ STO basis sets and double- ζ STO basis sets, respectively. The cores (N, O, C: 1s; S, Cl, P: 2p; W: 4d; Cu: 3p; Mo, Se, Br: 3d) were kept frozen. The Foster-Boys⁵ scheme has been applied to form the localized molecular orbital (LMOs). All calculations were performed by the ADF 2000 package⁶.

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1: M=W, X1=X2=X3=X4=S, X5=Cl 2: M=Mo, X1=O, X2=X3=X4=S, X5=Cl 3: M=Mo, X1=X2=X3=X4=S, X5=Se 4: M=Mo, X1=O, X2=X3=X4=S, X5=Br

The Mulliken charge density distribution of cluster **1**, WCu₃S₄Cl(PH₃)₃, is shown in **Figure 2**. The charge transfer (CT) occurs in two ways: (i) inside the metal core from Cu, S, Cl atoms to W atom; (ii) from metal core to the ligands (PH₃)₃ (MLCT). Considering the μ -bond effect, the first way of CT is dominate, which means the metal-to-metal CT (MMCT) in this type of transition metal cluster contributes to the optical nonlinearity.





Table 1 listed the composition percentages of the frontier orbitals of cluster **1**. We found that the HOMO and HOMO-1 are mainly consisted of the orbitals of Cu, S and Cl atoms, while those of P atoms only contribute a little (2.04% of HOMO). The orbitals of W and μ -S atoms dominated the LUMO and LUMO+1 orbitals. Thus we found again that the dominant CT inside the metal core from the bottom of Cu, S₄ and Cl to the top containing W.

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HOMO-1			НОМО			LUMO			LUMO+1		
ε (eV.)	-	4.674	ε (eV.)	-4	.645	ε (eV.)	-2.168		ε (eV.)	-2.150	
Cu1	3d	25.40	Cu1	3d	23.10	W	5d	47.13	W	5d	47.70
Cu2	3d	5.40		4p	1.20	S2	3p	17.40	S2	3p	11.90
Cu3	3d	4.25	Cu3	3d	9.24	S 3	3p	11.90	S4	3p	7.60
S 3	3p	20.30	Cu2	3d	2.30	S 8	3p	11.21	S 3	3p	7.20
S1	3p	8.20	Cl	3p	16.60	S1	3p	9.50	S1	3p	6.14
S2	3p	5.90	S2	3p	13.80	Cu	3d	5.03	Cu2	3d	3.75
S 4	3p	3.91	S 3	3p	7.80				Cu1	3d	1.38
Cl	3p	15.00	S4	3p	7.40						
			S 1	3p	7.30						
			P3	3p	2.04						

 Table 1
 Calculated percent population of frontier molecular orbitals of cluster 1

Table 2 shows the calculated β values and dipole moment of cluster **1–4**. Due to the symmetry of the cubane-like core structure, the cluster models have small β values comparing to those of the clusters with incomplete cubane-like structures. The largest β component is in *z* direction, $\beta_{av,z}$, because of the relatively poor structural symmetry of the metal core in this direction. The order of dipole moment is opposite to the order of β value. Compare cluster **1** to cluster **2**, the dipole increasing from 7.3 to 8.9 Debye while the β value is decreased from 4.1 to 1.2×10^{-30} esu. With the atom radius increases, the electronegativity decreases and the capability to form conjugation increases in the increase of the corresponding β values. This trend is different from the organic system, in which the β values depend on the ability of "push-pull" electrons.

Table 2 The calculated β (unit: 10⁻³⁰ esu) and dipole moment μ (unit: Debye) of cluster 1-4.

	1	2	3	4
$\beta_{av,x}$	0.2	0.4	0.2	0.7
$\beta_{av,y}$	2.4	0.8	0.3	1.3
$\beta_{av,z}$	3.4	0.8	2.1	3.2
β_{tot}	4.1	1.2	2.1	3.5
μ	7.3	8.6	6.5	6.2

In summary, the theoretical studies revealed that the CT inside the metal core of these series of cubane-like tetra-nuclear metal clusters contributes to the NLO property. The small first-order hyperpolarizabilities are due to their symmetric structures. The enhanced NLO property of this type of clusters could be obtained by substituting the ligands to result in the asymmetric cubane-like configurations.

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References

1. R. J. Sa, K. C. Wu, C. S. Lin, Chinese J. Struct. Chem., 2002, 21 (3), 336.

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- 2. a. A. Muller, H. Bogge. U. Schimanski, Inorg. Chim. Acta., 1983, 69, 5.
 - b. Q. F. Zhang, M. C. H, W. P. Su, R. Cao, H. Q. Liu, Polyhedron, 1997, 16, 1433.
 - c. N. Y. Zhu, R. C. Wu, X. T. Wu, Acta. Cryst.C., 1991, 47, 1537.
 - d. C. C. Lin, Z. S. Zheng, S. X. Liu, Chinese J. Struct. Chem., 1991, 10, 209.
 - e. Q. M. Wang, X. T. Wu, Q. Huang, T.L. Sheng, P. Lin Acta. Cryst.C., 1997, 53, 1542.
- 3. R. van Leeuwen., E. J. Baerends, Phys. Rev. A., 1994, 49, 2421.
- a. A.D. Becke, *Phys. Rev. A.*, **1988**, *38*, 3098.
 b. Lee, C., W. Yang, and R.G. Parr, *Phys. Rev. B.*, **1988**, *37* (2), 785.
- 5. S. F. Boys, Rev. Mod. Phys., 1960, 32, 296.
- 6. (a) E. J. Baerends, D. E. Ellis, A. Rosa. Chem. Phys., 1973, 2, 41. (b) G. te Velde, E. J. Baerends, J. Comput. Phys., 1992, 99, 84.

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